A splitting integrator for third order problems with transparent boundary conditions

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Introduction



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Water waves in shallow water

- Russell 1834: I was observing the motion of a boat which was rapidly drawn along a narrow channel by a pair of horses [...]
- Rayleigh and Boussinesq (1870), Korteweg and De Vries (1895): theoretical investigations
- Zabusky and Kruskal (1965): solitions
- Gardner, Greene, Kruskal and Miura (1967): inverse scattering transform



Figure. Russell's experiment (1995), Department of Mathematics, Heriot-Watt University, Edinburgh, Scotland

Collision of solitions https://www.youtube.com/ watch?v=wEbYELtGZwI

The Korteweg-de Vries (KdV) equation

$$\begin{cases} u_t + 6uu_x + u_{xxx} = 0, \quad t \in [0, T], \quad x \in \mathbb{R}, \\ u(0, x) = u^0(x). \end{cases}$$

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The Korteweg-de Vries (KdV) equation

$$\begin{cases} u_t + g(x)u_x + u_{xxx} = 0, & t \in [0, T], & x \in \mathbb{R}, \\ u(0, x) = u^0(x). \end{cases}$$

Motivation

Equation (2) finds applications in

- simulations of long ocean waves over an uneven bottom,
- propagation of fairly long waves in the shallow water.

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Numerical simulations require

- Temporal discretization
- Spatial discretization

(a)

(1)

The problem

Shallow Water Wave Properties



$$u_t + g(x)u_x + u_{xxx} = 0$$

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The problem

Shallow Water Wave Properties



$$u_t + g(x)u_x + u_{xxx} = 0$$

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Dispersion



Figure. Drogogne river, France. Picture: https://www.sudouest.fr/

$$u_t + u_{xxx} = 0$$

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Dispersion

$$u_t + u_{xxx} = 0$$

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Advection

$$u_t + g(x)u_x = 0$$

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Temporal discretization

Let m > 0 and $\tau = T/m$. Consider the *uniform* time discretization

$$0 = t^{0} < t^{1} < \dots < t^{m} = T, \quad t^{k} = k\tau.$$

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$$\begin{cases} u_{t} + g(x)u_{x} + u_{xxx} = 0, \\ u(0, x) = u^{0}(x), \end{cases}$$

$$\mathcal{T}^{t}(v^{0}):\begin{cases} v_{t}+g(x) v_{x}=0, \\ v(0,x)=v^{0}(x), \end{cases} \qquad \mathcal{D}^{t}(w^{0}):\begin{cases} w_{t}+w_{xxx}=0, \\ w(0,x)=w^{0}(x). \end{cases}$$

By Lie-Trotter splitting

$$u(\tau, x) \approx u^1(x) = \mathcal{D}^{\tau} \circ \mathcal{T}^{\tau}(u^0)$$

Boundary conditions



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Boundary conditions

Consider the problem

$$\begin{cases} u_t + g(x)u_x + u_{xxx} = 0, \quad t \in [0, T], \quad x \in \mathbb{R}, \\ u(0, x) = u^0(x). \end{cases}$$

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Boundary conditions

Consider the problem

$$\begin{cases} u_t + g(x)u_x + u_{xxx} = 0, & t \in [0, T], & x \in (a, b), \\ u(0, x) = u^0(x), & \\ u(t, a) = \dots & \\ u(t, b) = \dots & \\ u_x(t, b) = \dots & \\ \end{array}$$

... we must impose boundary conditions on *a* and *b*.

A typical choice are periodic boundary conditions.

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Periodic boundary conditions



Figure. Time simulation for T = 2 with periodic boundary conditions.

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Transparent boundary conditions



Figure. Time simulation for T = 2 with transparent boundary conditions.

Image: Image:

Transparent boundary conditions



Figure. Time simulation for T = 2 with transparent boundary conditions.

No free lunch theorem

- Transparent boundary conditions are difficult to compute
- Transparent boundary conditions are expensive to compute
- Transparent boundary conditions must be carefully design on the underlying numerical scheme

Transparent boundary conditions

We finally obtain the boundary conditions for the interior problem

$$\begin{cases} u^{m+1} + \tau u_{xxx}^{m+1} = u^m - \tau g(x)u_x^m, & m \ge 0, \quad x \in (a, b), \\ u(0, x) = u^0(x), \\ u^{m+1}(a) - Y_1^0 u_x^{m+1}(a) - Y_2^0 u_{xx}^{m+1}(a) = h_1^m, \\ u^{m+1}(b) - Y_3^0 u_{xx}^{m+1}(b) = h_2^m, \\ u_x^{m+1}(b) - Y_4^0 u_{xx}^{m+1}(b) = h_3^m. \end{cases}$$

where

$$h_1^m = \sum_{k=1}^{m+1} Y_1^k u_x^{m+1-k}(a) + Y_2^k u_{xx}^{m+1-k}(a),$$

$$h_2^m = \sum_{k=1}^{m+1} Y_3^k u_{xx}^{m+1-k}(b), \quad h_3^m = \sum_{k=1}^{m+1} Y_4^k u_{xx}^{m+1-k}(b).$$

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Given the semi-discrete scheme

$$\begin{cases} u^{m+1} + \tau u_{xxx}^{m+1} = u^m - \tau g(x) u_x^m & m \ge 0, \quad x \in [a, b], \\ u(0, x) = u^0(x), \\ u^{m+1}(a) - Y_1^0 u_x^{m+1}(a) - Y_2^0 u_{xx}^{m+1}(a) = h_1^m, \\ u^{m+1}(b) - Y_3^0 u_{xx}^{m+1}(b) = h_2^m, \\ u_x^{m+1}(b) - Y_4^0 u_{xx}^{m+1}(b) = h_3^m, \end{cases}$$

we perform a space discretization by a **pseudo-spectral method**. In particular, we implement a **dual-Petrov–Galerkin method**.

Numerical results



Figure. Left: time convergence, m number of time steps. Right: space convergence, N number of grid points.

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Numerical results



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Conclusions

- A splitting scheme was used for the time discretization
- Discrete transparent boundary conditions were designed for the particular time scheme
- A pseudo-spectral spatial discretization was presented to achieve fast convergence

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Future work

- Second order time scheme
- Rigorous stability analysis
- 2-D implementation

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Thanks for the attention!

L. Einkemmer, A. Ostermann, M. Residori, A pseudo-spectral splitting method for linear dispersive problems with transparent boundary conditions (arXiv-2019)

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